

Comment on 'The exact location of partition function zeros, a new method for statistical mechanics'

This article has been downloaded from IOPscience. Please scroll down to see the full text article.

1987 J. Phys. A: Math. Gen. 20 511

(<http://iopscience.iop.org/0305-4470/20/2/037>)

View [the table of contents for this issue](#), or go to the [journal homepage](#) for more

Download details:

IP Address: 129.252.86.83

The article was downloaded on 01/06/2010 at 05:21

Please note that [terms and conditions apply](#).

COMMENT

Comment on 'The exact location of partition function zeros, a new method for statistical mechanics'

A J Guttmann†

Department of Mathematics, University of Melbourne, Parkville, Victoria 3052, Australia

Received 12 May 1986

Abstract. In a recent paper Wood conjectured that the exact location of the branch point singularity for the hard square lattice gas on the negative activity z axis is at $z = -0.119\,392\dots$. A careful analysis of the available series data implies that this conjecture is almost certainly false. We find the correct value to be $z = -0.119\,338\,8809 \pm 0.000\,000\,001$, with corresponding exponent indistinguishably close to $\frac{5}{6}$ (corresponding to a zero of the singular part of the function).

Wood (1985) recently proposed a new mathematical technique for the precise location of sections of the locus of the limiting distribution of partition function zeros for model statistical mechanical systems. Motivated by certain properties of the Onsager solution of the two-dimensional Ising model, the conjecture was made for the hard square lattice gas that the branch point singularity of the activity series on the negative real axis is given by the closest negative root of the polynomial

$$2z^6 + 9z^5 + 42z^4 + 90z^3 + 96z^2 + 27z + 2$$

which yields as the root $z = -0.119\,392\dots$.

In order to investigate this conjecture, we have used the 42 term low-density series given by Baxter *et al* (1980) for the partition function per site κ as a function of the activity z . The coefficients of this series alternate in sign, suggesting that the dominant singularity is indeed on the negative real axis. Indeed, so dominant is this singularity that Baxter *et al* found no evidence of the physical singularity at $z_c \approx 3.7962$ from this series. In order to determine the nature of the closest singularity, we have utilised seven distinct methods of series analysis. We extrapolated the sequence of ratios, linear extrapolants of successive ratios, unbiased and biased exponent estimates by six sequence extrapolation methods. These were (i) Neville-Aitken extrapolation (Hartree 1952), (ii) the Barber-Hamer algorithm (Barber and Hamer 1982), (iii) Lubkin's three term formula (Lubkin 1952), (iv) Levin's u transform (Levin 1973), (v) the ε transformation (Wynn 1956) and (vi) Brezinski's θ algorithm (Brezinski 1971). We also utilised the method of integral approximants on the original series (Rehr *et al* 1980).

All these methods are described in Guttmann (1987), in which we argue that methods (i) and (vi) are the best sequence extrapolation methods for a wide variety of series in statistical mechanics, with method (iv) being nearly as good, while the

† Permanent address: Department of Mathematics, Statistics and Computer Science, University of Newcastle, Newcastle, NSW 2038, Australia.

Table 1. Estimates of the position of the singularity on the negative real axis of the low density $\kappa(z)$ series for the hard square lattice gas, with corresponding exponent estimates.

Method	Estimate of singularity position		Estimate of exponent	
	From ratio sequence	From linear extrapolants	Unbiased	Biased ($z = 0.119\ 338\ 881$)
Neville-Aitken	$ z < 0.119\ 339\ 6$	$ z > 0.119\ 338\ 3$	> 0.829	$> 0.831\ 1$
Barber-Hamer	$ z = 0.119\ 338\ 9$	$ z = 0.119\ 338\ 9$	$\approx 0.833\ 5$	$\approx 0.833\ 38$
Lubkin	$ z = 0.119\ 338\ 8$	—	—	—
Levin u	$ z > 0.119\ 338\ 1$	$ z \leq 0.119\ 338\ 86$	$\approx 0.833\ 7$	$< 0.834\ 0$
ϵ algorithm	$ z < 0.119\ 92$	$ z < 0.119\ 342$	> 0.819	$> 0.822\ 6$
θ transform	$ z = 0.119\ 338\ 88$	$ z = 0.119\ 338\ 881$	$\approx 0.833\ 37$	$\approx 0.833\ 38$

integral approximant method was shown to work well on a wider range of series than the sequence extrapolation algorithms.

In table 1 we show the results of the six sequence extrapolation methods. The implied uncertainty is in the last decimal digit quoted. It can be seen that all six methods are sufficiently precise to rule out Wood's conjecture. Further, they all agree amongst themselves. The most accurate method is Brezinski's θ transform, which gives

$$z^* = -0.119\ 338\ 8809 \pm 0.000\ 000\ 001$$

with an (unbiased) exponent estimate of 0.833 37, from which it is an obvious conjecture that the exponent is exactly $\frac{5}{6}$. This means that the low density series behaves like

$$\kappa(z) \sim (z - z^*)^{5/6} \quad z \rightarrow z^*.$$

The method of integral approximants, while less precise than the best of the sequence extrapolation methods, gives the estimates $z^* = -0.119\ 3390$, with exponent estimates greater than 0.831. Thus all seven methods are consistent, and all rule out the conjecture. It is interesting to note that Gaunt and Fisher (1965) were the first to consider this exponent, and noted that it was 'approximately 0.8'.

I would like to thank Professor R J Baxter for bringing the problem to my attention, and the ARGS for financial support.

References

Barber M N and Hamer C J 1982 *J. Aust. Math. Soc.* **13** 229
 Baxter R J, Enting I G and Tang S K 1980 *J. Stat. Phys.* **22** 465
 Brezinski C 1971 *C. R. Acad. Sci., Paris A* **273** 727
 Gaunt D S and Fisher M E 1965 *J. Chem. Phys.* **43** 2840
 Guttmann A J 1987 *Phase Transitions and Critical Phenomena* vol 11, ed C Domb and J L Lebowitz (New York: Academic) to appear
 Hartree D R 1952 *Numerical Analysis* (Oxford: Oxford University Press)
 Levin D 1973 *Int. J. Comput. Math.* **133** 371
 Lubkin S 1952 *J. Res. NBS B* **48** 228
 Rehr J J, Joyce G S and Guttmann A J 1980 *J. Phys. A: Math. Gen.* **13** 1587
 Wood D W 1985 *J. Phys. A: Math. Gen.* **18** L917
 Wynn P 1956 *J. Math. Phys.* **35** 318